Stability, Multiplicity, and Sunspots (deriving solutions to linearized system & Blanchard-Kahn conditions)

Wouter J. Den Haan London School of Economics

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Content

• A lot on sunspots

- Showing you what Dynare does (sort of) and Blanchard-Kahn conditions
 - and an even simpler way based on time iteration (an idea of Pontus Rendahl)

Introduction Sunspots

- What do we mean with non-unique solutions?
 - multiple solution versus multiple steady states
- What are sunspots?

Terminology

- Definitions are very clear
 - (use in practice can be sloppy)

Model:

$$H(p_{+1},p)=0$$

Solution:

$$p_{+1} = f(p)$$

multiple steady states; unique solution if initial p is given; (many solutions if no initial p is given



Multiple steady states & sometimes multiple solutions



 u_t

From Den Haan (2007)

Large sunspots (around 2000 at the peak)



Past Sun Spot Cycles



Sun spots even had a "Great Moderation"

Current cycle (another big one?)



9 / 48

Cute NASA video

https://www.youtube.com/watch?v=UD5VViT08ME

Sunspots in economics

- **Definition:** a solution is a sunspot solution if it depends on a stochastic variable that only appears *outside* the system. So not part of the model environment
- Model:

$$0 = \mathbb{E}H(p_{t+1}, p_t, d_{t+1}, d_t)$$

 d_t : exogenous random variable

Sunspots in economics (Cass & Shell 1983)

• Non-sunspot solution:

$$p_t = f(p_{t-1}, p_{t-2}, \cdots, d_t, d_{t-1}, \cdots)$$

• Sunspot:

$$p_t = f(p_{t-1}, p_{t-2}, \cdots, d_t, d_{t-1}, \cdots, s_t)$$

$$s_t : \text{ random variable with } \mathbb{E}[s_{t+1}] = 0$$

Origin of sunspots in economics

- William Stanley Jevons (1835-82) explored empirical relationship between sunspot activity (that is, the real thing!!!) and the price of corn.
- Fortunately, Jevons had some other contributions as well, such as "Jevons Paradox". His work is considered to be the start of mathematical economics.

Jevons Paradox

"It is wholly a confusion of ideas to suppose that the economical use of fuel is equivalent to a diminished consumption. The very contrary is the truth."

THE REBOUND EFFE

William Stanley Jevons

a British economist and logician.

Sunspots and science

Why are sunspots attractive?

- sunspots: s_t matters, just because agents believe this
 - self-fulfilling expectations don't seem that unreasonable
- sunspots provide many sources of shocks
 - number of sizable fundamental shocks small

Sunspots and science

Why are sunspots not so attractive?

- Purpose of science is to come up with predictions
 - If there is one sunspot solution, there are zillion others as well
- Support for the conditions that make them happen not overwhelming
 - you need sufficiently large increasing returns to scale or externality

Obtaining linear solutions: Overview

Getting started

simple examples

2 General derivation of Blanchard-Kahn solution

- When unique solution?
- When multiple solution?
- When no (stable) solution?
- **3** When do sunspots occur?

•

Getting started

Model: $y_{t+1} = \rho y_t$

Getting started

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• infinite number of solutions, independent of the value of ho

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Getting started

Model:
$$\begin{array}{c} y_{t+1} = \rho y_t \\ y_0 \text{ is given} \end{array}$$

Getting started

Model:
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• unique solution, independent of the value of ho

Getting started

• Blanchard-Kahn conditions apply to models that add as a requirement that the series do not explode

 $y_{t+1} = \rho y_t$

Model:

 y_t cannot explode

- $\rho > 1$: unique solution, namely $y_t = 0$ for all t
- ho < 1: many solutions
- $\rho = 1$: many solutions
 - be careful with ho=1, uncertainty matters

Neoclassical growth model; 2nd-order difference equation

$$\begin{array}{l} (k_{t-1}^{\alpha} + (1-\delta)k_{t-1} - k_t)^{-\gamma} \\ = \\ \beta(k_t^{\alpha} + (1-\delta)k_t - k_{t+1})^{-\gamma}(\alpha k_t^{\alpha-1} + 1 - \delta) \end{array}$$

 k_1 given

State-space representation

$$Ay_{t+1} + By_t = \varepsilon_{t+1}$$

 $\mathbb{E} [\varepsilon_{t+1} | I_t] = 0$
 $y_t: m \le n \text{ elements are not determined}$

some elements of ε_{t+1} are not exogenous shocks but prediction errors

Neoclassical growth model and state space representation

$$\mathbb{E} \begin{bmatrix} (\exp(z_{t})k_{t-1}^{\alpha} + (1-\delta)k_{t-1} - k_{t})^{-\gamma} = \\ \beta (\exp(z_{t+1})k_{t}^{\alpha} + (1-\delta)k_{t} - k_{t+1})^{-\gamma} \\ \times (\alpha \exp(z_{t+1})k_{t}^{\alpha-1} + 1 - \delta) \end{bmatrix} I_{t}$$

or equivalently without $\mathbb{E}\left[\cdot\right]$

$$(\exp(z_t)k_{t-1}^{\alpha} + (1-\delta)k_{t-1} - k_t)^{-\gamma} = \\ \beta (\exp(z_{t+1})k_t^{\alpha} + (1-\delta)k_t - k_{t+1})^{-\gamma} \\ \times (\alpha \exp(z_{t+1})k_t^{\alpha-1} + 1 - \delta) \\ + e_{\mathsf{E},t+1}$$

Neoclassical growth model and state space representation

Linearized model:

$$k_{t+1} = a_1k_t + a_2k_{t-1} + a_3z_{t+1} + a_4z_t + e_{\mathsf{E},t+1}$$

 $z_{t+1} = \rho z_t + e_{z,t+1}$
 k_0 is given

- k_t is end-of-period t capital
 - \implies k_t is chosen in t

Neoclassical growth model and state space representation

$$\begin{bmatrix} 1 & 0 & -a_3 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} k_{t+1} \\ k_t \\ z_{t+1} \end{bmatrix} + \begin{bmatrix} -a_1 & -a_2 & -a_4 \\ -1 & 0 & 0 \\ 0 & 0 & -\rho \end{bmatrix} \begin{bmatrix} k_t \\ k_{t-1} \\ z_t \end{bmatrix} = \begin{bmatrix} e_{\mathsf{E},t+1} \\ 0 \\ e_{z,t+1} \end{bmatrix}$$

Dynamics of the state-space system

$$Ay_{t+1} + By_t = \varepsilon_{t+1}$$

$$y_{t+1} = -A^{-1}By_t + A^{-1}\varepsilon_{t+1}$$
$$= Dy_t + A^{-1}\varepsilon_{t+1}$$

Thus

$$y_{t+1} = D^t y_1 + \sum_{l=1}^t D^{t-l} A^{-1} \varepsilon_{l+1}$$

Jordan matrix decomposition

$$D = P\Lambda P^{-1}$$

- Λ is a diagonal matrix with the eigen values of D
- without loss of generality assume that $|\lambda_1| \geq |\lambda_2| \geq \cdots |\lambda_n|$

Let

$$P^{-1} = \begin{bmatrix} \tilde{p}_1 \\ \vdots \\ \tilde{p}_n \end{bmatrix}$$

where \tilde{p}_i is a $(1 \times n)$ vector

Dynamics of the state-space system

$$y_{t+1} = D^{t}y_{1} + \sum_{l=1}^{t} D^{t-l}A^{-1}\varepsilon_{l+1}$$
$$= P\Lambda^{t}P^{-1}y_{1} + \sum_{l=1}^{t} P\Lambda^{t-l}P^{-1}A^{-1}\varepsilon_{l+1}$$

Dynamics of the state-space system

multiplying dynamic state-space system with P^{-1} gives

$$P^{-1}y_{t+1} = \Lambda^t P^{-1}y_1 + \sum_{l=1}^t \Lambda^{t-l} P^{-1} A^{-1} \varepsilon_{l+1}$$

or

$$\tilde{p}_i y_{t+1} = \lambda_i^t \tilde{p}_i y_1 + \sum_{l=1}^t \lambda_i^{t-l} \tilde{p}_i A^{-1} \varepsilon_{l+1}$$

recall that y_t is $n \times 1$ and \tilde{p}_i is $1 \times n$. Thus, $\tilde{p}_i y_t$ is a scalar

Model

$$\tilde{p}_i y_{t+1} = \lambda_i^t \tilde{p}_i y_1 + \sum_{l=1}^t \lambda_i^{t-l} \tilde{p}_i A^{-1} \varepsilon_{l+1}$$

$$\mathbb{E} \left[\varepsilon_{t+1} | I_t \right] = 0$$

- **3** m elements of y_1 are not determined
- **4** y_t cannot explode

Reasons for multiplicity

- **①** There are free elements in y_1
- **2** The only constraint on $e_{E,t+1}$ is that it is a prediction error.
 - This leaves lots of freedom

Eigen values and multiplicity

- Suppose that $|\lambda_1| > 1$
- To avoid explosive behavior it *must* be the case that

$$\begin{array}{l} \ \, \tilde{p}_1 y_1 = 0 \quad \text{and} \\ \ \, \boldsymbol{\hat{p}}_1 A^{-1} \varepsilon_l = 0 \quad \forall l \end{array}$$

How to think about #1?

$$\tilde{p}_1 y_1 = 0$$

- Simply an additional equation to pin down some of the free elements
- Much better: This is the policy function in the first period

How to think about #1?

$$\tilde{p}_1 y_1 = 0$$

Neoclassical growth model:

- $y_1 = [k_1, k_0, z_1]^T$
- $|\lambda_1|>1$, $|\lambda_2|<1$, $\lambda_3=
 ho<1$
- $\tilde{p}_1 y_1$ pins down k_1 as a function of k_0 and z_1
 - this is the policy function in the first period

How to think about #2?

$$\tilde{p}_1 A^{-1} \varepsilon_l = 0 \ \forall l$$

- This pins down $e_{E,t}$ as a function of $\varepsilon_{z,t}$
- That is, the prediction error must be a function of the structural shock, ε_{z,t}, and cannot be a function of other shocks,
 - i.e., there are no sunspots

How to think about #2?

$$\tilde{p}_1 A^{-1} \varepsilon_l = 0 \quad \forall l$$

Neoclassical growth model:

*p*₁A⁻¹ε_t says that the prediction error e_{E,t} of period t is a fixed function of the innovation in period t of the exogenous process, e_{z,t}

How to think about #1 combined with #2?

If these conditions on the RHS are imposed, then we get for the LHS

$$\tilde{p}_1 y_t = 0 \quad \forall t$$

- Without sunspots
 - i.e. with $\tilde{p}_1 A^{-1} \varepsilon_t = 0 \quad \forall t$
- k_t is pinned down by k_{t-1} and z_t in every period.

Blanchard-Kahn conditions

- Uniqueness: For every free element in y_1 , you need one $\lambda_i > 1$
- Multiplicity: Not enough eigenvalues larger than one
- No stable solution: Too many eigenvalues larger than one

How come this is so simple?

• In practice, it is easy to get

$$Ay_{t+1} + By_t = \varepsilon_{t+1}$$

• How about the next step?

$$y_{t+1} = -A^{-1}By_t + A^{-1}\varepsilon_{t+1}$$

- **Bad news**: *A* is often not invertible
- Good news: Same set of results can be derived
 - Schur decomposition (See Klein 2000 and Soderlind 1999)

Solutions to linear systems

- The analysis outlined above (requires A to be invertible)
- Generalized version of analysis above (see Klein 2000)
- Apply time iteration to linearized system (I learned this from Pontus Rendahl)

Solutions to linear systems

Model:

$$\Gamma_2 k_{t+1} + \Gamma_1 k_t + \Gamma_0 k_{t-1} = 0$$

or

$$\begin{bmatrix} \Gamma_2 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} k_{t+1} \\ k_t \end{bmatrix} + \begin{bmatrix} \Gamma_1 & \Gamma_0 \\ -1 & 0 \end{bmatrix} \begin{bmatrix} k_t \\ k_{t-1} \end{bmatrix} = 0$$

Standard approach #1

The method outlined above \Longrightarrow a unique solution of the form

$$k_t = ak_{t-1}$$

if BK conditions are satisfied

Standard approach #2

• Impose that the solution is of the recursive form

$$k_t = ak_{t-1}$$

and solve for a from

$$\Gamma_2 a^2 k_{t-1} + \Gamma_1 a k_{t-1} + \Gamma_0 k_{t-1} = 0 \quad \forall k_{t-1}$$

- Two solutions for $a: 0 < a_1 < 1, a_2 > 1$
- Does not simply generalize to higher-dimensional case

Time iteration

• Impose that the solution is of the form

 $k_t = ak_{t-1}$

- Use time iteration scheme, starting with $a_{[1]}$
- Recall that time iteration means using the guess for *tomorrows* behavior and then solve for *today*s behavior
- Method is demonstrated for scalar case but does easily generalize

(This simple procedure was pointed out to me by Pontus Rendahl)

Time iteration

- Follow the following iteration scheme, starting with $a_{[1]}$
 - Use $a_{[i]}$ to describe next period's behavior. That is,

$$\Gamma_2 a_{[i]} k_t + \Gamma_1 k_t + \Gamma_0 k_{t-1} = 0$$

(note the difference with last approach on previous slide)

• Obtain $a_{[i+1]}$ from

$$(\Gamma_2 a_{[i]} + \Gamma_1)k_t + \Gamma_0 k_{t-1} = 0$$

$$k_t = -\left(\Gamma_2 a_{[i]} + \Gamma_1\right)^{-1} \Gamma_0 k_{t-1}$$

$$a_{[i+1]} = -\left(\Gamma_2 a_{[i]} + \Gamma_1\right)^{-1} \Gamma_0$$

Advantages of time iteration

- It is simple, even if the "A matrix" is not invertible. (the inversion required by time iteration seems less problematic in practice)
- Since time iteration is linked to value function iteration, it has nice convergence properties

Example

$$k_{t+1} - 2k_t + 0.75k_{t-1} = 0$$

• The two solutions are

$$k_t = 0.5k_{t-1} \& k_t = 1.5k_{t-1}$$

 Time iteration on k_t = a_[i]k_{t-1} converges to stable solution for all initial values of a_[i] except 1.5.

References

- Larry Christiano taught me (a long time ago) this simple way of deriving the BK conditions and I think that I did not even change the notation.
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 - in case you want to do the analysis without the simplifying assumption that A is invertible
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 - also doesn't assume that A is invertible; possibly a more accessible paper